

Preliminary Thermal – August 2024

Easy Problems:

A1. For a wire with equation of state $F = bT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right)$, where L is the length, $L_0(T)$ is the length when the tension F is zero, b is a constant. Calculate the work done by the environment when the length changes from L_0 to $\frac{L_0}{2}$ at constant temperature.

Solution:

$$\begin{aligned} W &= \int F dL = \int_{L_0}^{\frac{L_0}{2}} bT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right) dL \\ &= bT \left(\frac{L^2}{2L_0} + \frac{L_0^2}{L} \right) \Bigg|_{L_0}^{\frac{L_0}{2}} = \frac{5}{8} bT L_0 \end{aligned}$$

A2. At low temperature, the constant volume molar specific heat of a solid is $C_v = a \left(\frac{T}{\theta_D} \right)^3$, where θ_D is called Debye temperature. Consider a solid with $a = 1.94$ kJ/mol/K and $\theta_D = 27$ °C, calculate the heat per mole the solid needs to absorb to increase temperature from 5 K to 10 K.

Solution:

$$Q = \int n C_v dT = \int_5^{10} 1.94 \times 10^3 \left(\frac{T}{300} \right)^3 dT = 0.168 J.$$

A3. Calculate $C_P = \left(\frac{\partial H}{\partial T} \right)_P$ using the chain rule for a simple solid for which $H(T, V) = Mc_0 T + (a_0 T - b_0)V$ and $V = V_0 \exp[(a_0 T - P)/b_0]$. Show that the result is $C_P = Mc_0 + \frac{a_0^2}{b_0} TV$.

Solution:

Using the chain rule, we have

$$C_P = \left(\frac{\partial H}{\partial T} \right)_P = \left(\frac{\partial H}{\partial T} \right)_V + \left(\frac{\partial H}{\partial V} \right)_T \left(\frac{\partial V}{\partial T} \right)_P$$

$$\left(\frac{\partial H}{\partial T} \right)_V = Mc_0 + a_0 V$$

$$\left(\frac{\partial H}{\partial V} \right)_T = a_0 T - b_0$$

$$\left(\frac{\partial V}{\partial T} \right)_P = \frac{a_0}{b_0} V$$

Therefore,

$$C_P = Mc_0 + a_0V + a_0T \frac{a_0}{b_0}V - b_0 \frac{a_0}{b_0}V = Mc_0 + \frac{a_0^2}{b_0}TV$$

Problem 5: (Easy) key
the state equation of ideal gas is

$$PV = nRT$$

The equation of adiabatic process is

$$P \left(\frac{V}{V_0} \right)^\gamma = P_0$$

$$\underline{P_0^{\text{He}} = P_0^{\text{air}}}$$

where $\gamma = C_p/C_v$, P_0 and P are starting and final pressures, respectively, and V_0 and V are volumes. Because $V_0 > V$ and $\gamma_{\text{He}} > \gamma_{\text{air}}$ one gets:

$$\frac{T_{\text{He}}}{T_{\text{air}}} = \frac{P_{\text{He}}}{P_{\text{air}}} = \frac{\left(\frac{V_0}{V} \right)^{\gamma_{\text{He}}}}{\left(\frac{V_0}{V} \right)^{\gamma_{\text{air}}}} = \left(\frac{V_0}{V} \right)^{\gamma_{\text{He}} - \gamma_{\text{air}}}$$

$$\frac{T_{\text{He}}}{T_{\text{air}}} = \left(\frac{V_0}{V} \right)^{\frac{7}{5} - \frac{5}{3}} = \left(\frac{V_0}{V} \right)^{\frac{6}{15}} > 1$$

Since $V_0 > V$. $\Rightarrow \underline{\underline{T_{\text{He}} > T_{\text{air}}}}$

Hard Problems:

B1. For ideal gas, if the heat capacity of a process is constant, then the process is polytropic $PV^l = \text{constant}$. Assuming that C_p and C_v are constants, find l .

Solution:

$$dU = C_n dT - PdV$$

$$(C_v - C_n)dT = -PdV = -nRT \frac{dV}{V}$$

$$(C_v - C_n) \frac{dT}{T} = -nR \frac{dV}{V}$$

$$(C_v - C_n) \ln T = -nR \ln V + \text{constant}$$

$$\ln(T^{C_v - C_n}) + \ln(V^{nR}) = \text{constant}$$

$$\exp[\ln(T^{C_v - C_n}) + \ln(V^{nR})] = \text{constant}$$

$$T^{C_v - C_n} V^{nR} = \text{constant}$$

$$(PV/nR)^{C_v - C_n} V^{nR} = \text{constant}$$

$$P^{C_v - C_n} V^{C_v - C_n + nR} = \text{constant}$$

$$P^{C_v - C_n} V^{C_p - C_n} = \text{constant}$$

$$P V^l = \text{constant}$$

$$l = (C_p - C_n) / (C_v - C_n)$$

B2. In the Einstein model of specific heat for diatomic molecules, atomic vibrations can be treated as independent quantum oscillators with quantized energy: $\mathcal{E}_n = n\hbar\omega$ and the population of excited oscillators follow the Boltzmann distribution.

- a) At thermal equilibrium, show that the average energy for the quantum oscillators is given by:

$$\bar{\mathcal{E}} = \frac{\hbar\omega}{e^{\hbar\omega/k_B T} - 1}.$$
- b) Find expression for the specific heat C_v of 1 mole of the diatomic molecules. Show that it approaches the $3R$ limit at high temperatures, where $R = N_A k_B$ is the universal gas constant.

Solution:

- a) For the Boltzmann distribution $e^{-\mathcal{E}_n/k_B T}$

$$\bar{\mathcal{E}} = \frac{\sum_{n=0}^{\infty} \mathcal{E}_n e^{-\mathcal{E}_n/k_B T}}{\sum_{n=0}^{\infty} e^{-\mathcal{E}_n/k_B T}} = -\frac{\partial}{\partial \beta} \ln[\sum_{n=0}^{\infty} e^{-\mathcal{E}_n \beta}] \quad \text{with } \beta = \frac{1}{k_B T}$$

$$\sum_{n=0}^{\infty} e^{-n\hbar\omega\beta} = \frac{1}{1 - e^{-\hbar\omega\beta}}$$

$$\bar{\mathcal{E}} = \frac{\partial}{\partial \beta} \ln[1 - e^{-\hbar\omega\beta}] = \frac{\hbar\omega}{e^{\hbar\omega/k_B T} - 1}$$

b) For 1 mole of the molecules,

$$U = 3N_A \frac{\hbar\omega}{e^{\hbar\omega/k_B T} - 1}$$

$$C_v = \left(\frac{\partial U}{\partial T} \right)_v = 3N_A k_B \left(\frac{\hbar\omega}{k_B T} \right)^2 \frac{e^{\hbar\omega/k_B T}}{(e^{\hbar\omega/k_B T} - 1)^2}$$

At high temperatures, $k_B T \gg \hbar\omega$

$$e^{\hbar\omega/k_B T} \approx 1 + \hbar\omega/k_B T$$

$$U = 3N_A k_B T$$

$$C_v = 3N_A k_B = 3R$$

B3. A solid with heat capacity C_A at temperature T_A is placed in contact with another solid with heat capacity C_B at a lower temperature T_B . What is the change in entropy of the system after the two bodies have reached thermal equilibrium?

Solution:

Assuming that the final temperature is T , according to energy conservation,

$$Q_A + Q_B = 0$$

$$\text{Or } C_A(T - T_A) + C_B(T - T_B) = 0$$

$$\text{So } T = (C_A T_A + C_B T_B) / (C_A + C_B).$$

In general, $dS = dQ/T = C/T dT$

$$\Delta S = \int dS = \int C d \ln T = \ln(T_f/T_i).$$

For the two systems:

$$\Delta S_A = C_A \ln(T/T_A).$$

$$\Delta S_B = C_B \ln(T/T_B).$$

$$\Delta S = \Delta S_A + \Delta S_B = C_A \ln(T/T_A) + C_B \ln(T/T_B)$$

$$= C_A \ln[(1 + C_B T_B / C_A T_A) / (1 + C_B / C_A)] + C_B \ln[(1 + C_A T_A / C_B T_B) / (1 + C_A / C_B)]$$

Problem 6 : (Hard) key

a) We consider nitrogen to be an ideal gas
The heat required is:

$$Q = n (C_v + R) \Delta T$$

$$= \frac{1000}{28} (5 + 2) \times 120 = \underline{\underline{30 \times 10^3 \text{ cal}}}$$

b) The increase of internal energy is

$$\Delta U = n C_v \Delta T = \frac{1000}{28} \times 5 \times 120$$

$$= 21.4 \times 10^3 \text{ cal}$$

c) The external work done is:

$$W = Q - \Delta U = 8.6 \times 10^3 \text{ cal}$$

d) If it is a process of constant volume, the required heat is

$$Q = n C_v \Delta T = 21.4 \times 10^3 \text{ cal}$$

CM

Peter 1

(A1) 5

$$\frac{\Delta L}{\Delta t} = \tau$$

5
(a)

$$\Delta L = \tau \Delta t = FR \Delta t$$

5

$$I\omega = FR \Delta t$$

5

$$\omega = \frac{FR \Delta t}{I} = \frac{18 \cdot 2.4 \cdot 15}{2100} = 3.09 \frac{\text{rad}}{\text{s}}$$

5 (6)

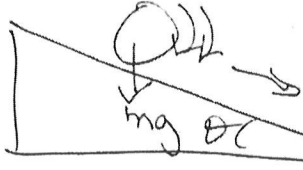
$$W = \frac{I\omega^2}{2} = \frac{2100 \cdot 3.09^2}{2} = 10^4 \text{ J} \approx 10 \text{ kJ}$$

CMA2

2) A hoop of radius 0.10 m and mass of 0.20 kg rolls down an inclined plane. If it starts from rest at a height of 2.0 m,

a) what is the total kinetic energy at the bottom of the incline (Use $g=9.8\text{m/s}^2$)?

7



$$mgh = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2 \quad I_{\text{hoop}} = mR^2$$

$$mgh \sin \theta = \frac{1}{2}mv^2 + \frac{1}{2}mR^2\omega^2$$

$$= 0.2\text{kg} \times 9.8\frac{\text{m}}{\text{s}^2} \times 2\text{m} = 3.92\text{J}$$

$\underbrace{\hspace{1.5cm}}_m \quad \underbrace{\hspace{1.5cm}}_g \quad \underbrace{\hspace{1.5cm}}_h$

b) what is the rotational kinetic energy at the bottom of the incline (the moment of inertia for a hoop is mR^2) ?

6

$$mgh = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2$$

$$= \frac{1}{2}mv^2 + \frac{1}{2}mR^2\omega^2$$

$$= \frac{1}{2}mv^2 + \frac{1}{2}mv^2$$

$$\frac{1}{2}I\omega^2 = \frac{1}{2}mgh = \frac{1}{2}3.92\text{J} = \underline{\underline{1.96\text{J}}}$$

c) what is the translational kinetic energy at the bottom of the incline ?

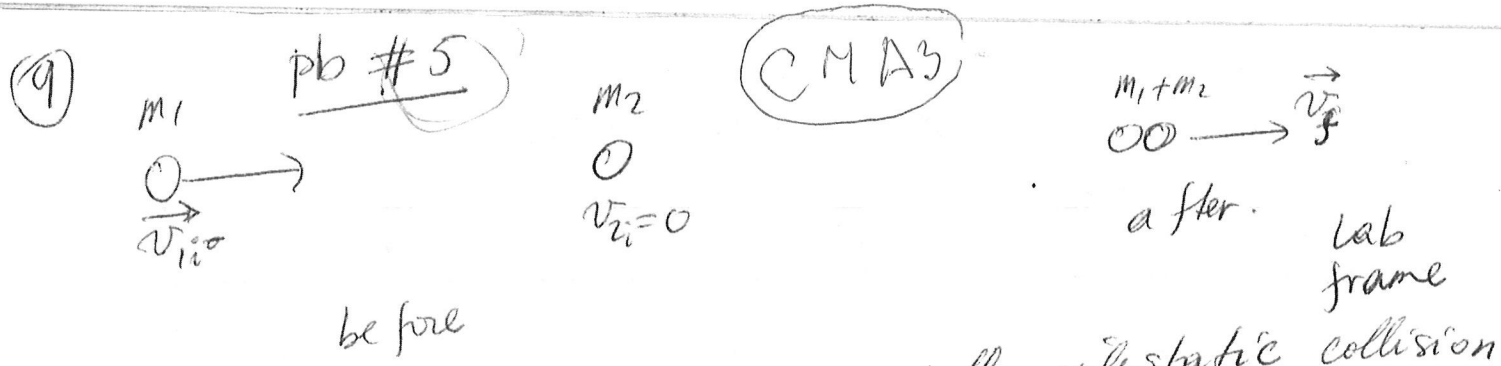
6

$$\frac{1}{2}mv^2 = mgh - \frac{1}{2}I\omega^2$$

$$\frac{1}{2}mv^2 = \frac{1}{2}mgh = \frac{1}{2}3.92\text{J} = \underline{\underline{1.96\text{J}}}$$

6

(d) $\omega = \sqrt{\frac{2E_{\text{rot}}}{I}} = \sqrt{\frac{2 \cdot 1.96}{0.2 \cdot 0.01}} = 44.3 \frac{\text{rad}}{\text{s}}$



a) momentum conservation for the perfectly inelastic collision

$$m_1 \vec{v}_{ii} + 0 = (m_1 + m_2) \vec{v}_f, \quad \vec{v}_f = \vec{v}_{if} = \vec{v}_{2f}$$

$$\Rightarrow \vec{v}_f = \frac{m_1}{m_1 + m_2} \vec{v}_{ii}$$

$$Q = T_i - T_f = \frac{1}{2} m_1 v_{ii}^2 - \frac{1}{2} (m_1 + m_2) v_f^2$$

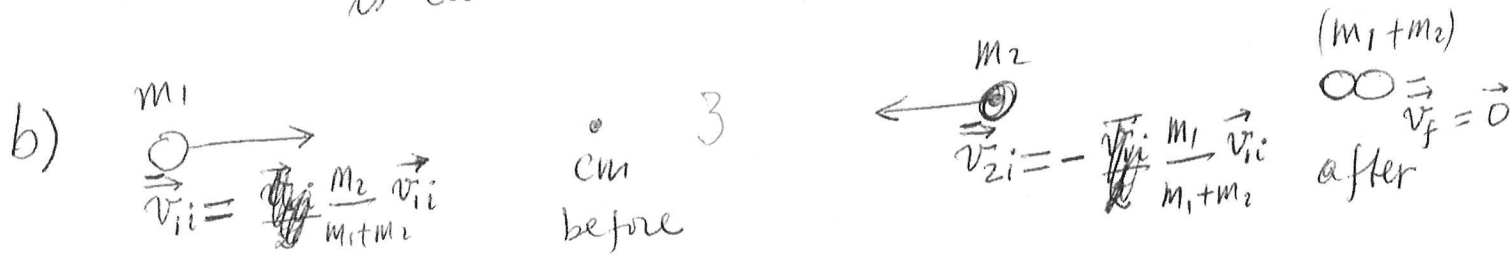
$$= \frac{1}{2} m_1 v_{ii}^2 - \frac{1}{2} (m_1 + m_2) \frac{m_1^2}{(m_1 + m_2)^2} v_{ii}^2$$

$$= \frac{1}{2} m_1 v_{ii}^2 \left(1 - \frac{m_1}{m_1 + m_2} \right)$$

$$Q = \frac{m_2}{m_1 + m_2} T_i = \frac{m_1 m_2}{2(m_1 + m_2)} v_{ii}^2 = \frac{1}{2} \mu v_{ii}^2$$

Alternative method: $Q = \frac{1}{2} \mu v_{ii}^2 (1 - e^2)$, $e=0$ for perfectly inelastic

$Q > 0 \Rightarrow$ an energy loss does occur, thus this process is called an exoergic collision.



(10)

CM A3, p. 2

$$Q = \bar{T}_i - \bar{T}_f$$

$$\bar{T}_f = 0$$

key result

$$2 \quad = \frac{1}{2} m_1 \bar{v}_1^2 + \frac{1}{2} m_2 \bar{v}_2^2 - 0$$

$$= \frac{1}{2} m_1 \left(\frac{m_2}{m_1 + m_2} \right)^2 v_{1i}^2 + \frac{1}{2} m_2 \left(\frac{m_1}{m_1 + m_2} \right)^2 v_{1i}^2$$

$$\left\{ \begin{array}{l} m_1 \vec{v}_{1i} + m_2 \vec{v}_{2i} = \vec{0} \\ m_1 \vec{v}_{1f} + m_2 \vec{v}_{2f} = \vec{0} \end{array} \right.$$

$$5 \quad = \frac{1}{2} \frac{m_1 m_2}{(m_1 + m_2)^2} (m_2 + m_1) v_{1i}^2$$

$$\begin{aligned} \vec{v}_{1f} &= \vec{0} \\ (m_1 + m_2) \vec{v}_{1f} &= \vec{0} \\ &= \vec{v}_{2f} = \vec{v}_f \end{aligned}$$

$$Q = \frac{1}{2} \frac{m_1 m_2}{m_1 + m_2} v_{1i}^2$$

the same as in the lab frame.

⑥ problem #3 (CMA4)

a) $F_d = -b \dot{x}$, find $v(x)$.

$$-b \dot{x} = m \frac{d\dot{x}}{dt} = m \frac{d\dot{x}}{dx} \frac{dx}{dt} = m \dot{x} \frac{d\dot{x}}{dx}$$

$$\Rightarrow \int_{x_0}^x -\frac{b}{m} dx = \int_{v_0}^v d\dot{x}$$

$$\Rightarrow v - v_0 = -\frac{b}{m}(x - x_0)$$

At $t=0$, $x = x_0 = 0$
 ~~$v = v_0 = 0$~~

$$\Rightarrow \boxed{v = v_0 - \frac{b}{m}x}$$

when $v=0$, $\boxed{x = \frac{mv_0}{b}}$

8

(A4)

$$m\ddot{x} = -c\dot{x}^2 \rightarrow m \frac{d\dot{x}}{dx} \dot{x} = -c\dot{x}^2 \quad \text{which is finite.}$$

b) $F_d = -c \dot{x}^2$, find $v(x)$ and $x(t)$

$$-c \dot{x}^2 = m \dot{x} \frac{d\dot{x}}{dx} \Leftrightarrow \int_{v_0}^v \frac{d\dot{x}}{\dot{x}} = -\int_{x_0}^x \frac{c}{m} dx$$

$$\Rightarrow \ln\left(\frac{v}{v_0}\right) = -\frac{c}{m}x, \quad \underline{x_0 = 0}$$

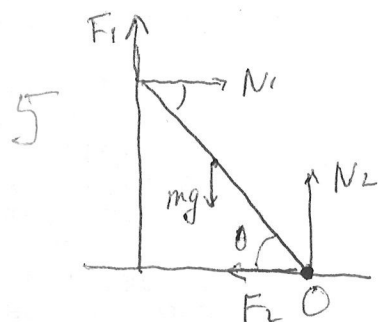
$$\Rightarrow \boxed{v(x) = v_0 e^{-\frac{c}{m}x}}$$

3

$v \rightarrow 0$ only when $x \rightarrow \infty$ (which is not finite), which means that the mass never comes to rest.

(B1)

CM



4

$$N_1 = F_2 \quad F_1 = \mu N_1$$

$$F_1 + N_2 = mg \quad F_2 = \mu N_2$$

↓

3

$$\left. \begin{aligned} N_1 &= \mu N_2 \\ \mu N_1 + N_2 &= mg \end{aligned} \right\} \begin{aligned} N_1 &= \mu (mg - \mu N_1) \\ N_1 &= \frac{\mu mg}{1 + \mu^2} \end{aligned}$$

equation for torques about point O:

3

$$N_1 l \sin \theta + \mu N_1 l \cos \theta - mg \frac{l}{2} \cos \theta = 0$$

$$\tan \theta = \frac{\frac{1}{2} mg - \mu N_1}{N_1} = \frac{mg}{2 N_1} - \mu$$

4

$$= \frac{mg(1 + \mu^2)}{2 \mu mg} - \mu = \frac{1}{2} \left(\frac{1}{\mu} + \mu \right) - \mu$$

$$= \frac{1}{2} \left(\frac{1}{\mu} - \mu \right)$$

2

Solve for μ $\mu^2 + 2\mu \tan \theta - 1 = 0$

$$\theta = 60^\circ$$

$$\cos \theta = \frac{1}{2} \quad \tan \theta = \sqrt{3}$$

$$\mu_{1,2} = -\tan \theta \pm \sqrt{\tan^2 \theta + 1} =$$

$$= -\tan \theta \pm \frac{1}{\cos \theta}$$

2

take positive root $\mu = -\sqrt{3} + 2 = 0.268$

(b)

2

$$F_2 = N_1 = \frac{\mu mg}{1 + \mu^2} = \frac{0.268 \cdot 70 \cdot 9.8}{1 + 0.268^2} = 172 \text{ N}$$

CM

(B2)

4) A string is wrapped around a hoop of radius 3 cm and mass 20 gms. The moment of inertia is mR^2

8 a) what is the angular acceleration in terms of $g (=9.8 \text{ m/s}^2)$?

7 $\tau = r \times F = I\alpha = TR \rightarrow T = \frac{I\alpha}{R}$

$mg - T = ma = m\alpha R$

$T = mg - ma = mg - m\alpha R$

$mg = m\alpha R + \frac{I\alpha}{R}$

$mg = \alpha(mR + \frac{I}{R})$

$\alpha = \frac{mg}{mR + \frac{I}{R}}$

$\alpha = \frac{g}{2R}$

[this is sufficient]

$\alpha = 163.3 \frac{\text{rad}}{\text{s}^2} = 16.7g$

5 b) calculate the time it takes for the hoop to descend 0.5 meters ?

$\alpha = \frac{g}{2R}$ $a = R\alpha$ so $a = \frac{g}{2} = 4.9 \frac{\text{m}}{\text{s}^2}$

$d = \frac{1}{2}at^2$

$\sqrt{\frac{2d}{a}} = t$ $t = 0.45 \text{ sec}$

5 c) Calculate the angular velocity of the rotating hoop after it has descended 0.5 m.

$\omega = \alpha t = 163.3 \frac{\text{rad}}{\text{s}^2} \cdot 0.45 \text{ sec} = 73.5 \frac{\text{rad}}{\text{sec}}$

$= 11.7 \frac{\text{rev.}}{\text{sec.}}$

$\frac{h}{2\pi R} = N = 2.65 \text{ rev}$

$\theta = 2\pi N = 16.7 \text{ rad}$

$\theta = \omega t + \frac{1}{2}\alpha t^2$

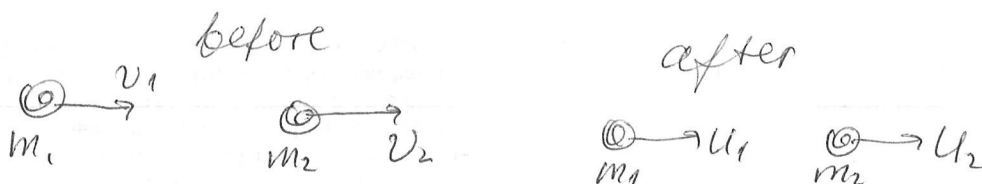
$\omega = \frac{\theta}{t} = \frac{16.7 \text{ rad}}{0.45 \text{ sec}} = 37.1 \frac{\text{rad}}{\text{sec}}$

$x = \frac{v_f^2 - v_i^2}{2a} \Rightarrow v_f = \sqrt{2ax}$

$v_f = R\omega$

CM

(B3)



$$6 \quad \begin{cases} \frac{m_1 v_1^2}{2} + \frac{m_2 v_2^2}{2} = \frac{m_1 u_1^2}{2} + \frac{m_2 u_2^2}{2} \\ m_1 v_1 + m_2 v_2 = m_1 u_1 + m_2 u_2 \end{cases}$$

$$3 \quad \begin{cases} m_1(v_1^2 - u_1^2) = m_2(u_2^2 - v_2^2) & (1) \\ m_1(v_1 - u_1) = m_2(u_2 - v_2) & (2) \end{cases}$$

3 1st solution: $v_1 = u_1, v_2 = u_2$
corresponds to no collisions

So assume $u_1 \neq v_1, u_2 \neq v_2$
then $v_1 + u_1 = u_2 + v_2$ divide (1) by (2)
and we have a linear system

2 mult. by m_1

$$u_1 - u_2 = v_2 - v_1$$

$$m_1 u_1 + m_2 u_2 = m_1 v_1 + m_2 v_2$$

$$m_2 u_2 + m_1 u_2 = m_1 v_1 + m_2 v_2 - m_1(v_2 - v_1)$$

$$u_2 = \frac{(m_2 - m_1)v_2 + 2m_1 v_1}{m_1 + m_2}$$

5

$$1 \rightleftharpoons 2 : u_1 = \frac{(m_1 - m_2)v_1 + 2m_2 v_2}{m_1 + m_2}$$

$$3 \quad (a) \quad m_1 = m_2 \quad u_2 = v_1 \quad u_1 = v_2$$

$$3 \quad (b) \quad m_1 \gg m_2 \quad u_2 = \frac{2v_1 - v_2}{2} = 2v_1 - v_2 \quad u_1 = 2v_2 - v_1$$

alternatively, one can use CM frame

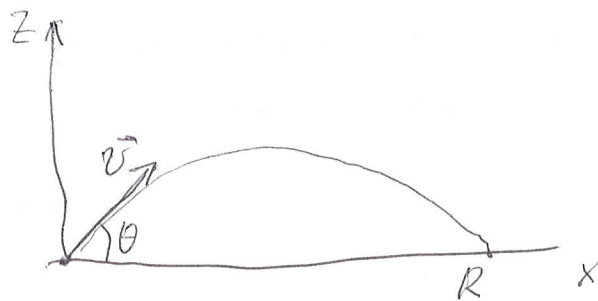
with $m_1 u_{1cm} = -m_2 u_{2cm}$ $|v_1 - v_2| = |u_{1cm} - u_{2cm}|$ relative velocities

and then add $u_1 = u_{1cm} + \frac{m_1 v_1 + m_2 v_2}{m_1 + m_2}$

$$u_{1cm} = \frac{m_2(v_2 - v_1)}{m_1 + m_2}$$

$$u_{2cm} = \frac{m_1(v_1 - v_2)}{m_1 + m_2}$$

CM (B4)
(a)



$$m\dot{v}_x = 0 \quad m\dot{v}_z = -g$$

2

$$x = v_{x0} t \quad z = v_{z0} t - \frac{g t^2}{2}$$

$$\text{for } x = R \quad z = 0 \rightarrow t = \frac{R}{v_{x0}}$$

3

$$0 = \frac{v_{z0}}{v_{x0}} R - \frac{g}{2} \left(\frac{R}{v_{x0}} \right)^2 = 0 \quad R = \frac{2 v_{x0} v_{z0}}{g} = \frac{2 v_0^2 \sin \theta \cos \theta}{g}$$

$$= \frac{v_0^2 \sin 2\theta}{g} = \frac{5^2 \cdot 10^4 \cdot \frac{\sqrt{3}}{2}}{9.8} = 2.21 \times 10^4 \text{ m} = 22.1 \text{ km}$$

$$(6) \quad m\dot{v}_x = -c_1 v_x, \quad m\dot{v}_z = -c_1 v_z - g$$

1

$$v_x = v_{x0} e^{-\gamma t}, \quad \gamma = \frac{c_1}{m} \quad v_z = \left(v_{z0} + \frac{g}{\gamma} \right) e^{-\gamma t} - \frac{g}{\gamma}$$

$$x = -\frac{v_{x0}}{\gamma} e^{-\gamma t} + \text{const} \quad x(0) = 0 \rightarrow \text{const} = \frac{v_{x0}}{\gamma}$$

2

$$x = \frac{v_{x0}}{\gamma} (1 - e^{-\gamma t}) \quad (1)$$

$$z = -\frac{1}{\gamma} \left(v_{z0} + \frac{g}{\gamma} \right) e^{-\gamma t} - \frac{g}{\gamma} t + \text{const}$$

$$z(0) = 0 \rightarrow \text{const} = \frac{1}{\gamma} \left(v_{z0} + \frac{g}{\gamma} \right)$$

2

$$z = \frac{1}{\gamma} \left(v_{z0} + \frac{g}{\gamma} \right) (1 - e^{-\gamma t}) - \frac{g}{\gamma} t \quad (2)$$

$$\text{for } x = R \quad z = 0$$

$$2 \quad \text{From (1)} \quad 1 - e^{-\gamma t} = \frac{R \gamma}{v_{x0}}, \quad t = -\frac{1}{\gamma} \ln \left(1 - \frac{R \gamma}{v_{x0}} \right)$$

Substitute into (2) with $z = 0$

CM (B4), p. 2

$$0 = \frac{1}{\gamma^1} \left(v_{z0} + \frac{g}{\gamma^1} \right) \frac{R \gamma^1}{v_{x0}} + \frac{g}{\gamma^{12}} \ln \left(1 - \frac{R \gamma^1}{v_{x0}} \right)$$

denote $u = \frac{R \gamma^1}{v_{x0}}$, then

$$2 \quad \left(v_{z0} + \frac{g}{\gamma^1} \right) u + \frac{g}{\gamma^1} \left(-u - \frac{u^2}{2} - \frac{u^3}{3} \right) = 0$$

where we used $\ln(1-u) = -u - \frac{u^2}{2} - \frac{u^3}{3}$

$$u v_{z0} - u^2 \frac{g}{2\gamma^1} - u^3 \frac{g}{3\gamma^1} = 0$$

$$5 \quad \frac{1}{3} u^3 + \frac{1}{2} u - v_{z0} \frac{\gamma^1}{g} = 0$$

$$u = \frac{-\frac{1}{2} + \sqrt{\frac{1}{4} + \frac{4}{3} v_{z0} \frac{\gamma^1}{g}}}{\frac{2}{3}} = \frac{3}{4} \left[-1 + \sqrt{1 + \frac{16}{3} v_{z0} \frac{\gamma^1}{g}} \right]$$

$$\gamma^1 = \frac{c_1}{m}$$

$$\gamma^1 = \frac{6.65 \times 10^{-5} \frac{\text{kg}}{\text{s}}}{0.05 \text{ kg}} = 1.33 \times 10^{-3} \text{ s}^{-1}$$

$$v_{x0} = 500 \frac{\text{m}}{\text{s}} \cdot \cos 30^\circ = 433 \frac{\text{m}}{\text{s}} \quad v_{z0} = 500 \frac{\text{m}}{\text{s}} \cdot \frac{1}{2} = 250 \frac{\text{m}}{\text{s}}$$

5

$$u = \frac{3}{4} \left(-1 + \sqrt{1 + \frac{16}{3} \cdot 250 \cdot \frac{1.33 \times 10^{-3}}{9.8}} \right) = 0.0650$$

$$R = u \frac{v_{x0}}{\gamma^1} = 0.065 \frac{433}{1.33 \times 10^{-3}} = 21161 \text{ m} = 21.16 \text{ km}$$